# Interaction of graphene with metallic and semiconductor surfaces. Ab initio approach to the lattice dynamics

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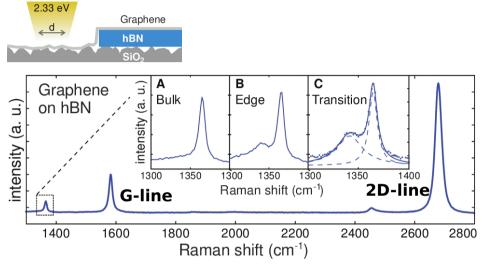


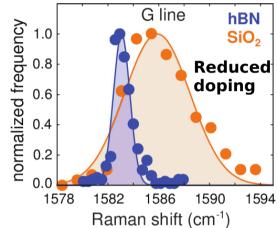
#### Interaction graphene/substrate

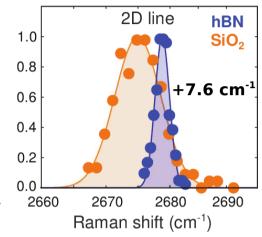
- Graphene has interesting properties
  - Test bench for fundamental physics.
- Response of graphene to the substrate interaction
  - Boron nitride, silicon carbide, iridium.
  - Raman and electron energy loss spectroscopy.
  - · Lattice dynamics. Density functional theory and GW
    - Attachment graphene/surface.
    - Influence of dielectric screening.
    - Persistence (or not) of some graphene fingerprints (Kohn anomaly, Dirac cone, etc).

## **Graphene@BN**

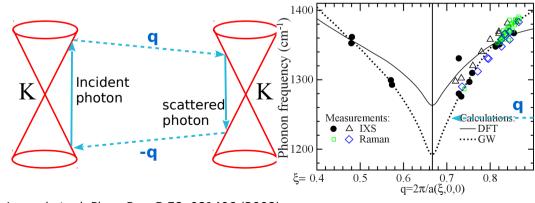
#### Boron nitride is one of the suitable substrates to keep intrinsic graphene



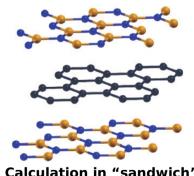




#### 2D-line comes from a 2-phonon process

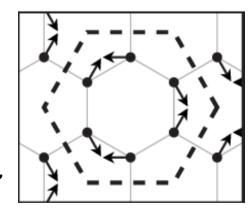


Lazzeri et. al. Phys. Rev. B 78, 081406 (2008) S. Berciaud et. al., Nano Lett. 10, 4074 (2010)



Calculation in "sandwich" graphene (conserving symmetry)

Slope of the TO phonon at K is proportional to the electron-phonon coupling.
GW approximation must be used.

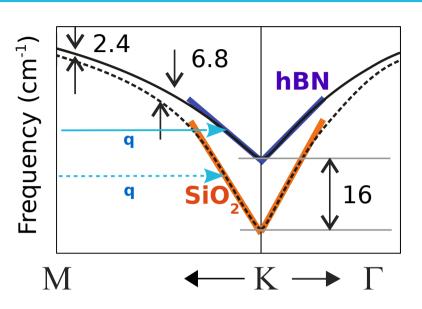


 $\langle D_{\mathbf{K}}^2 \rangle_F = \lim_{d \to 0} \frac{1}{8} \left( \frac{\Delta E_{\mathbf{K}}}{d} \right)^2$ 

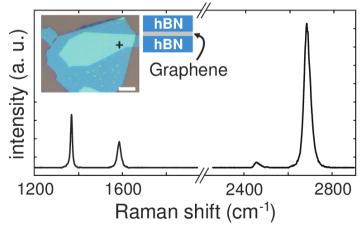
#### **Graphene@BN**

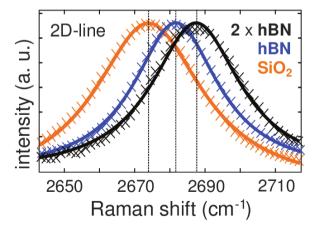
	graphene	on BN
$D_{\kappa}(LDA)$	89.25	86.00
$\overline{D_{\kappa}(GW)}$	207.88	191.27

**Dielectric screening** reduces the bandgap and the strength of the electron-phonon coupling.



- The slope of the optical phonon branch decreases.
- This also explains the down-shift in suspended graphene.
   Nano Lett. 1, 346 (2009).



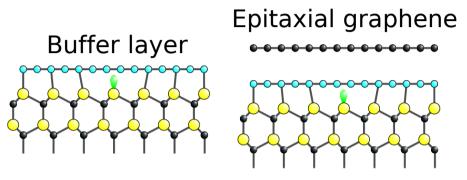


- Our approach can be tested by making a BN-graphene-BN sandwich.
- We observe double up-shift for the 2D-line.

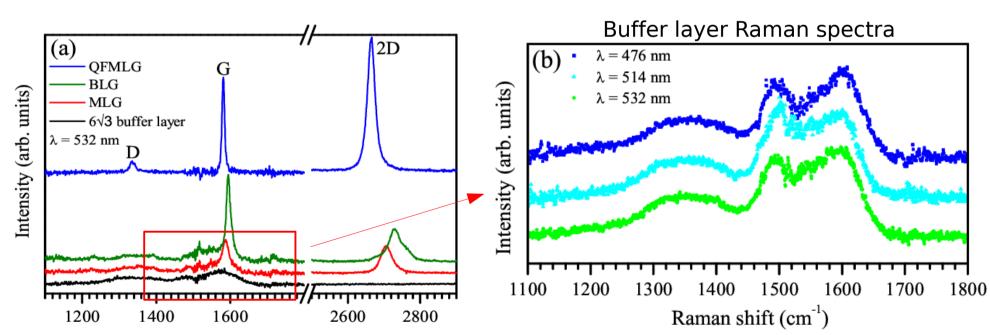
Dielectric screening is the responsible of the up-shift.

# Graphene@SiC. Buffer layer

Contribution of the buffer layer to the Raman spectrum from epitaxial graphene on SiC?

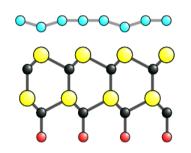


- Hybridization  $\pi$ -states and SiC states.
- Spectra are not compose of discrete peaks.
- Resemble a density of states?



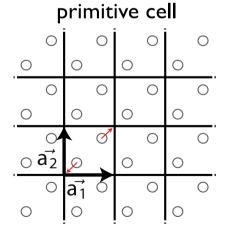
Substraction of SiC spectra

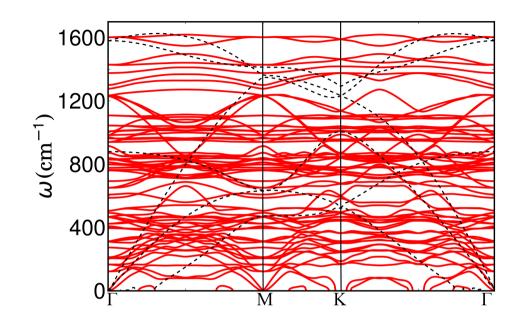
# Graphene@SiC. Buffer layer

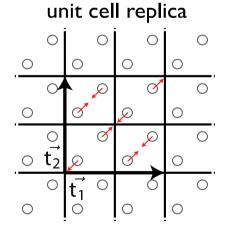


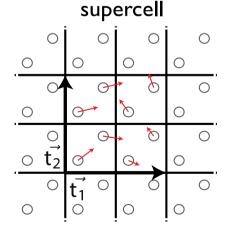
- Strong bonding C-Si: buckling of 0.04 nm
- Large cell to commensurate graphene and SiC
- Performing calculations in large supercells is like an origami, we fold the dispersion relation.
- We need to unfold the phonon modes to make easier the interpretation of our results.



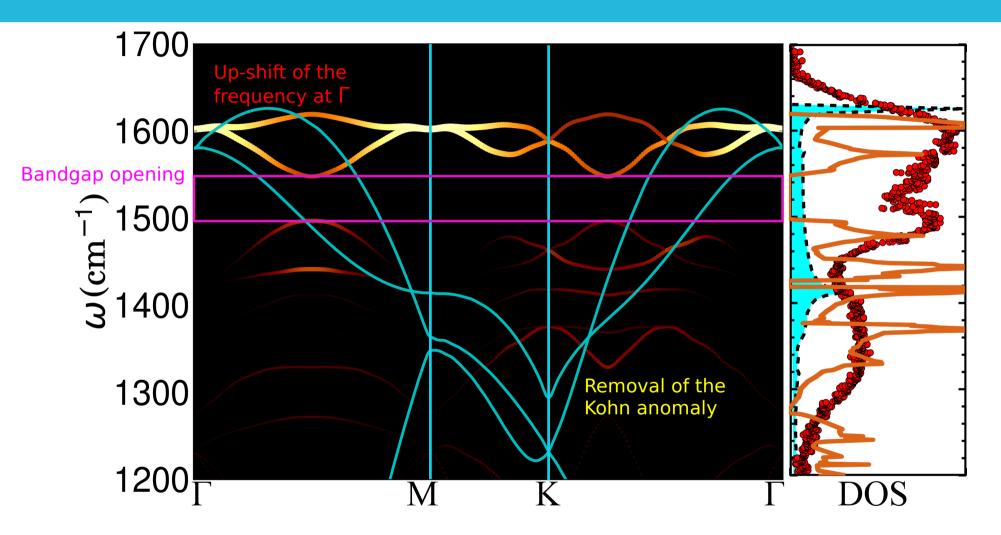








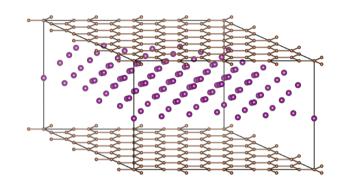
## Graphene@SiC. Buffer layer



DOS can be compared with Raman spectra

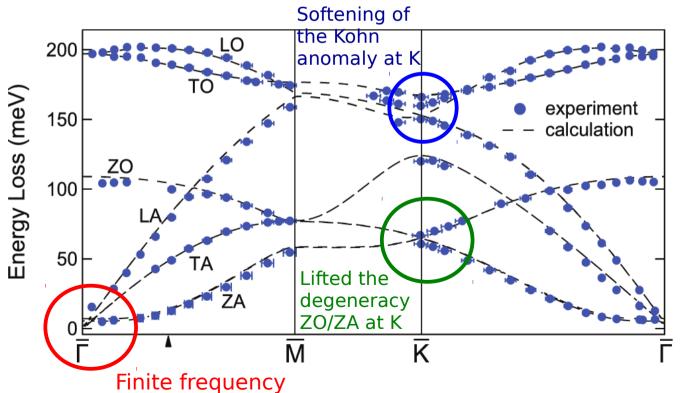
Different with respect to hBN

## Graphene@Iridium(111)



of ZA mode at Γ

- Graphene is very detached (d=3.5 nm)
- Lattice parameters are not commensurate.
   Ir(111) = 5.131Bohr, Graphene = 4.630 Bohr
- Formation of Moiré patterns
- Calculations in large supercells





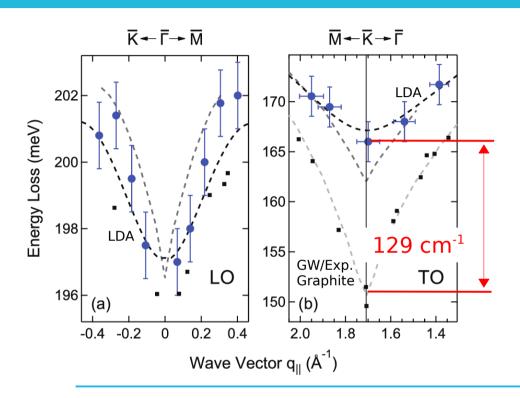
Diffraction spots: longrange Moiré pattern

Calculations performing with LDA.

Graphene unit cell, compressing iridium.

M. Endlich, AMS, L. Wirtz, J. Kröger. Phys. Rev. B 88, 205403 (2013)

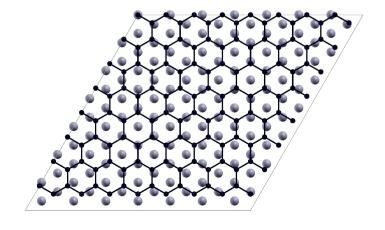
#### Graphene@Iridium(111)



- LDA calculations in graphene unit cell.
- The metallic screening cancels almost entirely the GW correction.

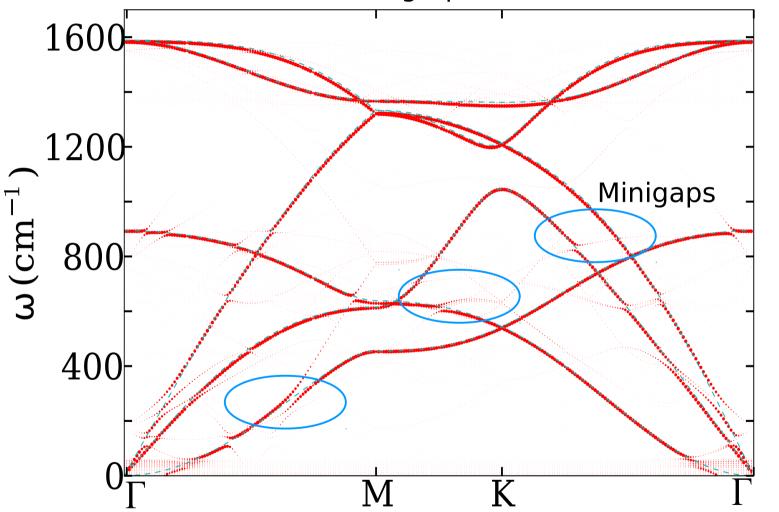
	graphene	on Ir	on BN
D <sub>K</sub> (LDA)	89.25	89.25	86.00
$D_{\kappa}(GW)$	207.88	131.75	191.27

- Variation of the local environment produces a corrugation in graphene.
- Number of atoms in unit cell exceeds the limit of application of ab-initio methods.
- Empirical methods. Force constant model.



# Graphene@Iridium(111)

Preliminar results for a 8x8-graphene unit cell (177 atoms)



Corrugation is still missing in the modeling Graphene phonon bands are almost unchanged

#### **Conclusions and future work**

 Lattice dynamics gives valuable information about attachment of graphene, screening, and conservation of intrinsic properties.

# Acknowledgements

- F. Fromm and T. Seyller (U. Chemnitz, DE).
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#### Conclusions and futur



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J. Kröger (TU Ilmenau, DE).



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