# Equivalent Harnack and gradient inequalities for pointwise curvature lower bound ${ }^{\$ 1}$ 

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#### Abstract

By using a coupling method, an explicit log-Harnack inequality with local geometry quantities is established for (sub-Markovian) diffusion semigroups on a Riemannian manifold (possibly with boundary). This inequality as well as the consequent $L^{2}$-gradient inequality, are proved to be equivalent to the pointwise curvature lower bound condition together with the convexity or absence of the boundary. Some applications of the log-Harnack inequality are also introduced.


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## 1. Introduction

Let $M$ be a $d$-dimensional connected complete Riemannian manifold possibly with a boundary $\partial M$. Consider $L=\Delta+Z$ for a $C^{1}$-vector field $Z$. Let $X_{t}(x)$ be the (reflecting) diffusion

[^0]process generated by $L$ with starting point $x$ and life time $\zeta(x)$. Then the associated diffusion semigroup $P_{t}$ is given by
$$
P_{t} f(x):=\mathbb{E}\left[f\left(X_{t}(x)\right) 1_{\{t<\zeta(x)\}}\right], \quad t \geqslant 0, f \in \mathscr{B}_{b}(M) .
$$

Although the semigroup depends on $Z$ and the geometry on the whole manifold, we aim to establish Harnack, resp. gradient type inequalities for $P_{t}$ by using local geometry quantities.

Let $K \in C(M)$ be such that

$$
\begin{equation*}
\operatorname{Ric}_{Z}:=\operatorname{Ric}-\nabla Z \geqslant-K, \tag{1.1}
\end{equation*}
$$

i.e. for any $x \in M$ and $X \in T_{x} M, \operatorname{Ric}(X, X)-\left\langle X, \nabla_{X} Z\right\rangle \geqslant-K(x)|X|^{2}$. Next, for any $D \subset M$, let

$$
K(D):=\sup _{D} K, \quad D_{r}=\{z \in M: \rho(z, D) \leqslant r\}, \quad r \geqslant 0,
$$

where $\rho$ is the Riemannian distance on $M$. Finally, to investigate $P_{t}$ using local curvature bounds, we introduce, for a given bounded open domain $D \subset M$, the following class of reference functions:

$$
\mathscr{C}_{D}=\left\{\phi \in C^{2}(\bar{D}):\left.\phi\right|_{D}>0,\left.\phi\right|_{\partial D \backslash \partial M}=0,\left.N \phi\right|_{\partial M \cap \partial D} \geqslant 0\right\},
$$

where $N$ is the inward unit normal vector field of $\partial M$. When $\partial M=\emptyset$, the restriction $\left.N \phi\right|_{\partial M} \geqslant 0$ is automatically dropped. For any $\phi \in \mathscr{C}_{D}$, we have

$$
c_{D}(\phi)=\sup _{D}\left\{5|\nabla \phi|^{2}-\phi L \phi\right\} \in[0, \infty) .
$$

The finiteness of $c_{D}(\phi)$ is trivial since $\bar{D}$ is compact. To see that $c_{D}(\phi) \geqslant 0$, we consider the following two situations:
(a) There exists $x \in \partial D \backslash \partial M$. We have $\phi(x)=0$ so that $c_{D}(\phi) \geqslant\left\{5|\nabla \phi|^{2}-\phi L \phi\right\}(x) \geqslant 0$.
(b) When $\partial D \backslash \partial M=\emptyset$, we have $\bar{D}=M$. Otherwise, there exists $z \in M \backslash(D \cup \partial M)$. For any $z^{\prime} \in D \backslash \partial M$, let $\gamma:[0,1] \rightarrow M \backslash \partial M$ be a smooth curve linking $z$ and $z^{\prime}$. Since $z^{\prime} \in D$ but $z \notin D$, there exists $s \in[0,1]$ such that $\gamma(s) \in \partial D$. This is however impossible since $\partial D \subset$ $\partial M$ and $\gamma(s) \notin \partial M$. Therefore, in this case $M=\bar{D}$ is compact so that the reflecting diffusion process is non-explosive. Now, let $x \in \bar{D}$ such that $\phi(x)=\max _{\bar{D}} \phi$. Since $\left.N \phi\right|_{\partial M} \geqslant 0$ due to $\phi \in \mathscr{C}_{D}, \phi\left(X_{t}\right)-\phi(x)-\int_{0}^{t} L \phi\left(X_{s}\right) \mathrm{d} s$ is a sub-martingale so that

$$
\phi(x) \geqslant \mathbb{E} \phi\left(X_{t}\right) \geqslant \phi(x)+\int_{0}^{t} \mathbb{E} L \phi\left(X_{s}\right) \mathrm{d} s, \quad t \geqslant 0 .
$$

This implies $L \phi(x) \leqslant 0$ (known as the maximum principle) and thus,

$$
c_{D}(\phi) \geqslant\left\{5|\nabla \phi|^{2}-\phi L \phi\right\}(x) \geqslant 0
$$

Theorem 1.1. Let $K \in C(M)$. The following statements are equivalent:
(1) (1.1) holds and $\partial M$ is either empty or convex.
(2) For any bounded open domain $D \subset M$ and any $\phi \in \mathscr{C}_{D}$, the log-Harnack inequality

$$
\begin{aligned}
& P_{T} \log f(y)-\log \left(P_{T} f(x)+1-P_{T} 1(x)\right) \\
& \quad \leqslant \frac{\rho(x, y)^{2}}{2}\left(\frac{K\left(D_{\rho(x, y)}\right)}{1-\mathrm{e}^{-2 K\left(D_{\rho(x, y)}\right) T}}+\frac{c_{D}(\phi)^{2}\left(\mathrm{e}^{2 K\left(D_{\rho(x, y)) T}\right.}-1\right)}{2 K\left(D_{\rho(x, y))}\right) \phi(y)^{4}}\right), \\
& \\
& T>0, y \in D, x \in M,
\end{aligned}
$$

holds for strictly positive $f \in \mathscr{B}_{b}(M)$.
(3) For any bounded open domain $D \subset M$ and any $\phi \in \mathscr{C}_{D}$,

$$
\left|\nabla P_{T} f\right|^{2}(x) \leqslant\left\{P_{T} f^{2}-\left(P_{T} f\right)^{2}\right\}(x)\left(\frac{K(D)}{1-\mathrm{e}^{-2 K(D) T}}+\frac{c_{D}(\phi)^{2}\left(\mathrm{e}^{2 K(D) T}-1\right)}{2 K(D) \phi(x)^{4}}\right)
$$

holds for all $x \in D, T>0, f \in \mathscr{B}_{b}(M)$.
If moreover $P_{T} 1=1$, then the statements above are also equivalent to:
(4) For any bounded open domain $D \subset M$ and any $\phi \in \mathscr{C}_{D}$, the Harnack type inequality

$$
\begin{aligned}
P_{T} f(y) \leqslant & P_{T} f(x) \\
& +\rho(x, y)\left(\frac{K(D)}{1-\mathrm{e}^{-2 K(D) T}}+\frac{c_{D}(\phi)^{2}\left(\mathrm{e}^{2 K(D) T}-1\right)}{2 K(D) \inf _{\ell(x, y)} \phi^{4}}\right)^{1 / 2} \sqrt{P_{T} f^{2}(y)}
\end{aligned}
$$

holds for non-negative $f \in \mathscr{B}_{b}(M), T>0$ and $x, y \in D$ such that the minimal geodesic $\ell(x, y)$ linking $x$ and $y$ is contained in $D$.

Remark. (i) When $K$ is constant, a number of equivalent semigroup inequalities are available for the curvature condition (1.1) together with the convexity or absence of the boundary, see $[8,10]$ and references within (see also [3,11] for equivalent semigroup inequalities of the curvaturedimension condition). When $\partial M$ is either empty or convex, the above result provides at the first time equivalent semigroup properties for the general pointwise curvature lower bound condition.
(ii) When the diffusion process is explosive, the appearance of $1-P_{T} 1$ in the log-Harnack inequality is essential. Indeed, without this term the inequality does not hold for e.g. $f \equiv 1$ provided $P_{T} 1<1$.
(iii) The following result shows that the constant $1 / 2$ involved in the log-Harnack inequality is sharp.

Proposition 1.2. Let $c>0$ be a constant. For any $x \in M$, strictly positive function $f$ with $|\nabla f|(x)>0$ and $\log f \in C_{0}^{2}(M)$, and any constants $C>0$, the inequality

$$
P_{T} \log f(y) \leqslant \log \left(P_{T} f(x)+1-P_{T} 1(x)\right)+c \rho(x, y)^{2}\left(\frac{C}{1-\mathrm{e}^{-2 C T}}+\mathrm{o}\left(\frac{1}{T}\right)\right)
$$

for small $T>0$ and small $\rho(x, y)$ implies that $c \geqslant 1 / 2$.
Proof. Let us take $v \in T_{x} M$ and $y_{s}=\exp _{x}[s v], s \geqslant 0$. Then the given log-Harnack inequality implies that

$$
\begin{equation*}
P_{s} \log f\left(y_{s}\right)-\log \left(P_{s} f(x)+1-P_{s} 1(x)\right) \leqslant c s^{2}|v|^{2}\left(\frac{C}{1-\mathrm{e}^{-2 C s}}+\mathrm{o}\left(\frac{1}{s}\right)\right) \tag{1.2}
\end{equation*}
$$

holds for small $s>0$. On the other hand, for any $g \in C^{2}(M)$ with bounded $L g$, one has

$$
\begin{equation*}
\left.\frac{\mathrm{d}}{\mathrm{~d} s} P_{s} g\right|_{s=0}=L g \tag{1.3}
\end{equation*}
$$

Indeed, letting $X_{t}$ be the diffusion process generated by $L$ with $X_{0}=x$, by Itô's formula and the dominated convergence theorem we obtain

$$
\lim _{s \downarrow 0} \frac{P_{s} g(x)-g(x)}{s}=\lim _{s \downarrow 0} \frac{1}{s} \mathbb{E} \int_{0}^{s \wedge \zeta(x)} \operatorname{Lg}\left(X_{r}\right) \mathrm{d} r=\mathbb{E} \lim _{s \downarrow 0} \frac{1}{s} \int_{0}^{s \wedge \zeta(x)} \operatorname{Lg}\left(X_{r}\right) \mathrm{d} r=L g(x) .
$$

Combining (1.2) with (1.3) we obtain

$$
\begin{aligned}
\langle v, \nabla \log f\rangle(x)-|\nabla \log f|^{2}(x) & =L \log f(x)+\langle v, \nabla \log f\rangle(x)-\frac{L f(x)}{f(x)} \\
& =\lim _{s \downarrow 0} \frac{1}{s}\left\{P_{s} \log f\left(y_{s}\right)-\log \left(P_{s} f(x)+1-P_{s} 1(x)\right)\right\} \\
& \leqslant \frac{c|v|^{2}}{2} .
\end{aligned}
$$

Taking $v=r \nabla \log f(x)$ for $r \geqslant 0$ we obtain

$$
\left(r-1-\frac{c r^{2}}{2}\right)|\nabla \log f(x)|^{2} \leqslant 0, \quad r \geqslant 0 .
$$

This implies $c \geqslant 1 / 2$ by taking $r=1 / c$.
To derive the explicit log-Harnack inequality using local geometry quantities, we may take e.g. $D=B(y, 1):=\{z: \rho(y, z)<1\}$. Let

$$
\begin{aligned}
& K_{y}=0 \vee K(B(y, 1)), \quad K_{x, y}=K(B(y, 1+\rho(x, y))), \\
& K_{y}^{0}=0 \vee \sup \left\{-\operatorname{Ric}(U, U): U \in T_{z} M,|U|=1, z \in B(y, 1)\right\}, \\
& b_{y}=\sup _{B(y, 1)}|Z| .
\end{aligned}
$$

Then $K\left(D_{\rho(x, y)}\right)=K_{x, y}$ and according to [7, proof of Corollary 5.1] (see page 121 therein with $\bar{\delta}_{x}$ replaced by 1 ), we may take $\phi(z)=\cos \frac{\pi \rho(y, z)}{2}$ so that $\phi(y)=1$ and

$$
\kappa(y):=K_{y}+\frac{\pi^{2}(d+3)}{4}+\pi\left(b_{y}+\frac{1}{2} \sqrt{K_{y}^{0}(d-1)}\right) \geqslant c_{D}(\phi) .
$$

Note that when $\partial M$ is convex, $\left.N \rho(\cdot, y)\right|_{\partial M} \leqslant 0$ so that $\left.N \phi\right|_{\partial D \cap \partial M} \geqslant 0$ as required in the definition of $\mathscr{C}_{D}$. Therefore, Theorem 1.1(2) implies that

$$
\begin{align*}
P_{t} \log f(y) \leqslant & \log \left\{P_{t} f(x)+1-P_{t} 1(x)\right\} \\
& +\frac{\rho(x, y)^{2}}{2}\left(\frac{K_{x, y}}{1-\mathrm{e}^{-2 K_{x, y} t}}+\frac{\kappa(y)^{2}\left(\mathrm{e}^{2 K_{x, y} t}-1\right)}{2 K_{x, y}}\right) \tag{1.4}
\end{align*}
$$

holds for all strictly positive $f \in \mathscr{B}_{b}(M), x, y \in M$ and $t>0$. As in the proofs of [6, Corollary 1.2] and [9, Corollary 1.3], this implies the following heat kernel estimates and entropy-cost inequality. When $P_{t}$ obeys the log-Sobolev inequality for $t>0$, the second inequality in Corollary 1.3(2) below also implies the HWI inequality as shown in [4,5].

Corollary 1.3. Assume (1.1) and that $\partial M$ is either convex or empty. Let $Z=\nabla V$ for some $V \in C^{2}(M)$ such that $P_{t}$ is symmetric w.r.t. $\mu(\mathrm{d} x):=\mathrm{e}^{V(x)} \mathrm{d} x$, where $\mathrm{d} x$ is the volume measure. Let $p_{t}$ be the density of $P_{t}$ w.r.t. $\mu$. Assume that (1.1) holds.
(1) Let $\bar{K}(y)=K(B(y, 2))$. Then

$$
\begin{aligned}
& \int_{M} p_{t}(y, z) \log p_{t}(y, z) \mu(\mathrm{d} z) \\
& \leqslant \sqrt{t \wedge 1}\left(\frac{\bar{K}(y)}{1-\mathrm{e}^{-2 \bar{K}(y) t}}+\frac{\kappa(y)^{2}\left(\mathrm{e}^{2 \bar{K}(y) t}-1\right)}{2 \bar{K}(y)}\right)+\log \frac{P_{2 t} 1(y)+\mu\left(1-P_{t} 1\right)}{\mu(B(y, \sqrt{t \wedge 1}))}
\end{aligned}
$$

holds for all $y \in M$ and $t>0$.
(2) If $\mu$ is a probability measure and $P_{t} 1=1$, then the Gaussian heat kernel lower bound

$$
\begin{aligned}
& p_{2 t}(x, y) \geqslant \exp \left[-\frac{\rho(x, y)^{2}}{2}\left(\frac{K_{x, y}}{1-\mathrm{e}^{-2 K_{x, y} t}}+\frac{\kappa(y)^{2}\left(\mathrm{e}^{2 K_{x, y} t}-1\right)}{2 K_{x, y}}\right)\right], \\
& \quad t>0, x, y \in M
\end{aligned}
$$

and the entropy-cost inequality

$$
\begin{aligned}
& \int_{M}\left(P_{t} f\right) \log P_{t} f \mathrm{~d} \mu \\
& \quad \leqslant \inf _{\pi \in \mathscr{C}(\mu, f \mu)} \int_{M \times M} \frac{\rho(x, y)^{2}}{2}\left(\frac{K_{x, y}}{1-\mathrm{e}^{-2 K_{x, y} t}}+\frac{\kappa(y)^{2}\left(\mathrm{e}^{2 K_{x, y} t}-1\right)}{2 K_{x, y}}\right) \pi(\mathrm{d} x, \mathrm{~d} y), \\
& t>0,
\end{aligned}
$$

hold for any probability density function $f$ of $\mu$, where $\mathscr{C}(\mu, f \mu)$ is the set of all couplings of $\mu$ and $f \mu$.

Proof. According to (1.4), the heat kernel lower bound in (2) follows from the proof of [9, Corollary 1.3], while the other two inequalities can be proved as in the proof of [6, Corollary 1.2]. Below we only present a brief proof of (1).

By an approximation argument we may apply (1.4) to $f(z):=p_{t}(y, z)$ so that

$$
\begin{aligned}
I & :=\int_{M} p_{t}(y, z) \log p_{t}(y, z) \mu(\mathrm{d} z) \\
& \leqslant \log \left\{p_{2 t}(x, y)+1-P_{t} 1(x)\right\}+\frac{\rho(x, y)^{2}}{2}\left(\frac{K_{x, y}}{1-\mathrm{e}^{-2 K_{x, y} t}}+\frac{\kappa(y)^{2}\left(\mathrm{e}^{2 K_{x, y} t}-1\right)}{2 K_{x, y}}\right) .
\end{aligned}
$$

Since $K_{x, y} \leqslant \bar{K}(y)$ for $x \in B(y, 1)$, this implies that

$$
\begin{aligned}
& \mathrm{e}^{I} \mu(B(y, \sqrt{t \wedge 1})) \exp \left[-\frac{t \wedge 1}{2}\left(\frac{\bar{K}(y)}{1-\mathrm{e}^{-2 \bar{K}(y) t}}+\frac{\kappa(y)^{2}\left(\mathrm{e}^{2 \bar{K}(y) t}-1\right)}{2 \bar{K}(y)}\right)\right] \\
& \leqslant \mathrm{e}^{I} \int_{M} \exp \left[-\frac{\rho(x, y)^{2}}{2}\left(\frac{K_{x, y}}{1-\mathrm{e}^{-2 K_{x, y} t}}+\frac{\kappa(y)^{2}\left(\mathrm{e}^{2 K_{x, y} t}-1\right)}{2 K_{x, y}}\right)\right] \mu(\mathrm{d} x)
\end{aligned}
$$

$$
\leqslant \int_{M}\left\{p_{2 t}(x, y)+1-P_{t} 1(x)\right\} \mu(\mathrm{d} x)=P_{2 t} 1(y)+\mu\left(1-P_{t} 1\right) .
$$

This proves (1).
We remark that the entropy upper bound in (1) is sharp for short time, since both $-\log \mu(B(y, \sqrt{t}))$ and the entropy of the Gaussian heat kernel behave like $\frac{d}{2} \log \frac{1}{t}$ for small $t>0$.

## 2. Proof of Theorem 1.1

We first observe that when $P_{T} 1=1$ the equivalence of (3) and (4) is implied by the proof of [12, Proposition 1.3]. Indeed, by (3)

$$
\left|\nabla P_{T} f\right|^{2} \leqslant\left\{P_{T} f^{2}-\left(P_{T} f\right)^{2}\right\}\left(\frac{K(D)}{1-\mathrm{e}^{-2 K(D) T}}+\frac{c_{D}(\phi)^{2}\left(\mathrm{e}^{2 K(D) T}-1\right)}{2 K(D) \inf _{\ell(x, y)} \phi^{4}}\right)
$$

holds on the minimal geodesic $\ell(x, y)$, so that the Harnack inequality in (4) follows from the first part in the proof of [12, Proposition 1.3]. On the other hand, by the second part of the proof, the inequality in (4) implies

$$
\left|\nabla P_{t} f\right|^{2} \leqslant\left\{P_{T} f^{2}\right\}\left(\frac{K(D)}{1-\mathrm{e}^{-2 K(D) T}}+\frac{c_{D}(\phi)^{2}\left(\mathrm{e}^{2 K(D) T}-1\right)}{2 K(D) \phi^{4}}\right)
$$

on $D$. Replacing $f$ by $f-P_{T} f(x)$, we obtain the inequality in (3) since $\nabla P_{T} f=\nabla P_{T}(f-$ $P_{T} f(x)$ ) provided $P_{T} 1=1$.

In the following three subsections, we prove (1) implying (2), (2) implying (3), and (3) implying (1) respectively.

### 2.1. Proof of (1) implying (2)

We assume the curvature condition (1.1) and that $\partial M$ is either empty or convex. To prove the log-Harnack inequality in (2), we will make use of the coupling argument proposed in [1]. As explained in [1, Section 3], we may and do assume that the cut-locus of the manifold is empty.

Now, let $T>0$ and $y \in D, x \neq y$ be fixed. For any $z, z^{\prime} \in M$, let $P_{z, z^{\prime}}: T_{z} M \rightarrow T_{z^{\prime}} M$ be the parallel transport along the unique minimal geodesic from $z$ to $z^{\prime}$. Let $X_{t}$ solve the following Itô type SDE on $M$

$$
\mathrm{d}^{I} X_{t}=\sqrt{2} \Phi_{t} \mathrm{~d} B_{t}+Z\left(X_{t}\right) \mathrm{d} t+N\left(X_{t}\right) \mathrm{d} l_{t}, \quad X_{0}=x
$$

up to the life time $\zeta(x)$, where $B_{t}$ is the $d$-dimensional Brownian motion, $\Phi_{t}$ is the horizontal lift of $X_{t}$ on the frame bundle $O(M)$, and $l_{t}$ is the local time of $X_{t}$ on $\partial M$ if $\partial M \neq \emptyset$. When $\partial M=\emptyset$, we simply take $l_{t}=0$ so that the last term in the equation disappears.

To construct another process starting at $y$ such that it meets $X_{t}$ before $T$ and its hitting time to $\partial D$, let $Y_{t}$ solve the SDE with $Y_{0}=y$

$$
\begin{equation*}
\mathrm{d}^{I} Y_{t}=\sqrt{2} P_{X_{t}, Y_{t}} \Phi_{t} \mathrm{~d} B_{t}+Z\left(Y_{t}\right) \mathrm{d} t-\sqrt{\xi_{1}(t)^{2}+\xi_{2}(t)^{2}} \nabla \rho\left(X_{t}, \cdot\right)\left(Y_{t}\right) \mathrm{d} t+N\left(Y_{t}\right) \mathrm{d} \tilde{l}_{t}, \tag{2.1}
\end{equation*}
$$

where $\tilde{l}_{t}$ is the local time of $Y_{t}$ on $\partial M$ when $\partial M \neq \emptyset$, and

$$
\begin{aligned}
& \xi_{1}(t)=\frac{2 K\left(D_{\rho(x, y)}\right) \exp \left[-K\left(D_{\rho(x, y)}\right) t\right]}{1-\exp \left[-2 K\left(D_{\rho(x, y)}\right) T\right]} \rho(x, y) 1_{\left\{Y_{t} \neq X_{t}\right\}}, \\
& \xi_{2}(t)=\frac{2 c_{D}(\phi) \rho\left(X_{t}, Y_{t}\right)}{\phi\left(Y_{t}\right)^{2}}, \quad t \in[0, T] .
\end{aligned}
$$

Then $Y_{t}$ is well-defined before $T \wedge \tau_{D(x, y)}(x) \wedge \tau_{D}(y)$, where

$$
\tau_{D}(y):=\inf \left\{t \in[0, T \wedge \zeta(x)): Y_{t} \in \partial D\right\}, \quad \tau_{D(x, y)}(x)=\inf \left\{t \geqslant 0: X_{t} \notin D(x, y)\right\} .
$$

Let

$$
\tau=\inf \left\{t \in[0, \zeta(x) \wedge \zeta(y)): X_{t}=Y_{t}\right\}
$$

where $\inf \emptyset=\infty$ by convention.
Let $\Theta=\tau \wedge T \wedge \tau_{D}(y) \wedge \tau_{D(x, y)}(x)$ and set

$$
\eta(t)=\frac{1}{\sqrt{2}} \sqrt{\xi_{1}(t)^{2}+\xi_{2}(t)^{2}} \nabla \rho\left(\cdot, Y_{t}\right)\left(X_{t}\right), \quad t \in[0, \Theta)
$$

Define

$$
R=\exp \left[-\int_{0}^{\Theta}\left\langle\eta(t), \Phi_{t} \mathrm{~d} B_{t}\right\rangle-\frac{1}{2} \int_{0}^{\Theta}|\eta(t)|^{2} \mathrm{~d} t\right]
$$

We intend to prove
(i) $R$ is a well-defined probability density with

$$
\mathbb{E}\{R \log R\} \leqslant \frac{\rho(x, y)^{2}}{2}\left(\frac{K\left(D_{\rho(x, y)}\right)}{1-\mathrm{e}^{-2 K\left(D_{\rho(x, y)}\right) T}}+\frac{c_{D}(\phi)^{2}\left(\mathrm{e}^{2 K\left(D_{\rho(x, y)}\right)^{T}}-1\right)}{2 K\left(D_{\rho(x, y)}\right) \phi(y)^{4}}\right)
$$

(ii) $\tau \leqslant T \wedge \tau_{D}(y) \wedge \tau_{D(x, y)}(x)$ holds $\mathbb{Q}$-a.s., where $\mathbb{Q}:=R \mathbb{P}$.

Once these two assertions are confirmed, by taking $Y_{t}=X_{t}$ for $t \geqslant \tau$ we see that $Y_{t}$ solves (2.1) up to its life time $\zeta(y)=\zeta(x)$ and $X_{T}=Y_{T}$ for $T<\zeta(x)$. Moreover, by the Girsanov theorem the process

$$
\tilde{B}_{t}:=B_{t}+\int_{0}^{t} \eta(s) \mathrm{d} s, \quad t \geqslant 0
$$

is a $d$-dimensional Brownian motion under $\mathbb{Q}$ and Eq. (2.1) can be reformulated as

$$
\begin{equation*}
\mathrm{d}^{I} Y_{t}=\sqrt{2} P_{X_{t}, Y_{t}} \Phi_{t} \mathrm{~d} \tilde{B}_{t}+Z\left(Y_{t}\right) \mathrm{d} t+N\left(Y_{t}\right) \mathrm{d} \tilde{l}_{t}, \quad Y_{0}=y . \tag{2.2}
\end{equation*}
$$

Combining this with the Young inequality (see [2, Lemma 2.4])

$$
\begin{aligned}
P_{T} \log f(y)= & \mathbb{E}\left\{R 1_{\{T<\zeta(y)\}} \log f\left(Y_{T}\right)\right\}=\mathbb{E}\left\{R 1_{\{T<\zeta(x)\}} \log f\left(X_{T}\right)\right\} \\
\leqslant & \mathbb{E} R \log R+\log \mathbb{E} \exp \left[1_{\{T<\zeta(x)\}} \log f\left(X_{T}\right)\right] \\
= & \log \left(P_{T} f(x)+1-P_{T} 1(x)\right)+\mathbb{E} R \log R \\
\leqslant & \log \left(P_{T} f(x)+1-P_{T} 1(x)\right) \\
& +\frac{\rho(x, y)^{2}}{2}\left(\frac{K\left(D_{\rho(x, y)}\right)}{1-\mathrm{e}^{-2 K\left(D_{\rho(x, y))} T\right.}-1}+\frac{c_{D}(\phi)^{2}\left(\mathrm{e}^{2 K\left(D_{\rho(x, y)}\right) T}-1\right)}{2 K\left(D_{\rho(x, y)}\right) \phi(y)^{4}}\right) .
\end{aligned}
$$

This gives the desired log-Harnack inequality.

Below we prove (i) and (ii) respectively.
Lemma 2.1. For any $n \geqslant 1$, let

$$
\tau_{n}(y)=\inf \left\{t \in[0, T \wedge \zeta(x)): \rho\left(Y_{t}, D^{c}\right) \leqslant n^{-1}\right\}
$$

and

$$
\Theta_{n}=\tau \wedge \frac{n T}{n+1} \wedge \tau_{D(x, y)}(x) \wedge \tau_{n}(y)
$$

Let $R_{n}$ be defined as $R$ using $\Theta_{n}$ in place of $\Theta$. Then $\left\{R_{n}\right\}_{n} \geqslant 1$ is a uniformly integrable martingale with $\mathbb{E} R_{n}=1$ and

$$
\mathbb{E}\left\{R_{n} \log R_{n}\right\} \leqslant \frac{\rho(x, y)^{2}}{2}\left(\frac{K\left(D_{\rho(x, y}\right)}{1-\mathrm{e}^{-2 K\left(D_{\rho(x, y)}\right) T}-1}+\frac{c_{D}(\phi)^{2}\left(\mathrm{e}^{2 K\left(D_{\rho(x, y)}\right) T}-1\right)}{2 K\left(D_{\rho(x, y)}\right) \phi(y)^{4}}\right)
$$

for $n \geqslant 1$. Consequently, (i) holds.
Proof. (i) follows from the first assertion and the martingale convergence theorem. Since before time $\Theta_{n}$ the process $\eta(t)$ is bounded, the martingale property and $\mathbb{E} R_{n}=1$ is well-known. So, it remains to prove the entropy upper bound. By the Itô formula we see that (cf. (2.3) and (2.4) in [1])

$$
\begin{equation*}
\mathrm{d} \rho\left(X_{t}, Y_{t}\right) \leqslant K\left(D_{\rho(x, y)}\right) \rho\left(X_{t}, Y_{t}\right) \mathrm{d} t-\sqrt{\xi_{1}(t)^{2}+\xi_{2}(t)^{2}} \mathrm{~d} t, \quad t \leqslant \Theta_{n} \tag{2.3}
\end{equation*}
$$

Then

$$
\mathrm{d} \rho\left(X_{t}, Y_{t}\right)^{2} \leqslant 2 K\left(D_{\rho(x, y)}\right) \rho\left(X_{t}, Y_{t}\right)^{2} \mathrm{~d} t-\frac{4 c_{D}(\phi) \rho\left(X_{t}, Y_{t}\right)^{2}}{\phi\left(Y_{t}\right)^{2}} \mathrm{~d} t, \quad t \leqslant \Theta_{n}
$$

Note that $\left(\tilde{B}_{t}\right)_{t \in\left[0, \Theta_{n}\right]}$ is a $d$-dimensional Brownian motion under the probability $\mathbb{Q}_{n}:=R_{n} \mathbb{P}$. Combining this with (2.2) and using Itô's formula along with the facts that the martingale part of $\rho\left(X_{t}, Y_{t}\right)^{2}$ is zero and $\left.N \phi\right|_{\partial D \cap \partial M} \geqslant 0$, we obtain

$$
\begin{aligned}
\mathrm{d}\left\{\frac{\rho\left(X_{t}, Y_{t}\right)^{2}}{\phi\left(Y_{t}\right)^{4}}\right\} \leqslant & \mathrm{d} M_{t}-\frac{4 \rho\left(X_{t}, Y_{t}\right)^{2}}{\phi\left(Y_{t}\right)^{6}}\left\{c_{D}(\phi)+\phi\left(Y_{t}\right) L \phi\left(Y_{t}\right)-5\left|\nabla \phi\left(Y_{t}\right)\right|^{2}\right\} \mathrm{d} t \\
& -\frac{2 K\left(D_{\rho(x, y)) \rho\left(X_{t}, Y_{t}\right)^{2}}^{\phi\left(Y_{t}\right)^{4}} \mathrm{~d} t\right.}{} \\
\leqslant & \mathrm{d} M_{t}-\frac{2 K\left(D_{\rho(x, y))}\right)\left(X_{t}, Y_{t}\right)^{2}}{\phi\left(Y_{t}\right)^{4}} \mathrm{~d} t, \quad t \leqslant \Theta_{n}
\end{aligned}
$$

where

$$
\mathrm{d} M_{t}:=-\frac{4 \rho\left(X_{t}, Y_{t}\right)^{2}}{\phi\left(Y_{t}\right)^{5}}\left\langle\nabla \phi\left(Y_{t}\right), P_{X_{t}, Y_{t}} \Phi_{t} \mathrm{~d} \tilde{B}_{t}\right\rangle
$$

is a $\mathbb{Q}_{n}$-martingale for $t \leqslant \Theta_{n}$. This implies

$$
\mathbb{E}_{\mathbb{Q}_{n}}\left\{\frac{\rho\left(X_{t \wedge \Theta_{n}}, Y_{t \wedge \Theta_{n}}\right)^{2}}{\phi\left(Y_{t \wedge \Theta_{n}}\right)^{4}}\right\} \leqslant \frac{\rho(x, y)^{2}}{\phi(y)^{4}} \mathrm{e}^{2 K\left(D_{\rho(x, y)) t}\right.}, \quad t \geqslant 0 .
$$

Hence,

$$
\begin{aligned}
\mathbb{E}\left\{R_{n} \log R_{n}\right\}= & \frac{1}{2} \mathbb{E}_{\mathbb{Q}_{n}} \int_{0}^{\Theta_{n}}|\eta(t)|^{2} \mathrm{~d} t=\frac{1}{4} \mathbb{E}_{\mathbb{Q}_{n}} \int_{0}^{\Theta_{n}}\left\{\xi_{1}(t)^{2}+\xi_{2}(t)^{2}\right\} \mathrm{d} t \\
\leqslant & \frac{K\left(D_{\rho(x, y)}\right)^{2} \rho(x, y)^{2}}{\left(1-\mathrm{e}^{-2 K\left(D_{\rho(x, y)}\right)}\right)^{2}} \int_{0}^{T} \mathrm{e}^{-2 K\left(D_{\rho(x, y))^{\prime} t} \mathrm{~d} t\right.} \\
& +c_{D}(\phi)^{2} \int_{0}^{T} \mathbb{E}_{\mathbb{Q}_{n}} \frac{\rho\left(X_{t \wedge \Theta_{n}}, Y_{\left.t \wedge \Theta_{n}\right)^{2}}\right.}{\phi\left(Y_{\left.t \wedge \Theta_{n}\right)^{4}}\right.} \mathrm{d} t \\
\leqslant & \frac{K\left(D_{\rho(x, y)}\right) \rho(x, y)^{2}}{2\left(1-\mathrm{e}^{-2 K\left(D_{\rho(x, y)}\right) T}\right)}+\frac{c_{D}(\phi)^{2}\left(\mathrm{e}^{2 K\left(D_{\rho(x, y)) T}\right.}-1\right) \rho(x, y)^{2}}{2 K\left(D_{\rho(x, y))}\right) \phi(y)^{4}}, \quad s>0 .
\end{aligned}
$$

Lemma 2.2. We have $\tau \leqslant T \wedge \tau_{D}(y) \wedge \tau_{D(x, y)}(x), \mathbb{Q}$-a.s.
Proof. By (2.3) we have

$$
\begin{equation*}
\int_{0}^{\Theta}\left\{\xi_{1}(t)+\xi_{2}(t)\right\} \mathrm{d} t=\lim _{n \rightarrow \infty} \int_{0}^{\Theta_{n}}\left\{\xi_{1}(t)+\xi_{2}(t)\right\} \mathrm{d} t<\infty \tag{2.4}
\end{equation*}
$$

Since under $\mathbb{Q}$ the process $Y_{t}$ is generated by $L$, as observed at the beginning of [7, Section 4] we have

$$
\int_{0}^{\tau_{D}(y)} \frac{1}{\Phi\left(Y_{t}\right)^{2}} \mathrm{~d} t=\infty, \quad \mathbb{Q} \text {-a.s. }
$$

Then (2.4) implies that $\mathbb{Q}$-a.s.

$$
\begin{equation*}
\tau_{D}(y)>\tau_{D(x, y)}(x) \wedge \tau \wedge T \tag{2.5}
\end{equation*}
$$

Moreover, it follows from (2.3) that

$$
\begin{aligned}
\rho\left(X_{t}, Y_{t}\right) & \leqslant \mathrm{e}^{K\left(D_{\rho(x, y))} t\right.} \rho(x, y)-\int_{0}^{t} \mathrm{e}^{K\left(D_{\rho(x, y)}\right)(t-s)} \xi_{1}(s) \mathrm{d} s \\
& \leqslant \frac{\mathrm{e}^{-2 K\left(D_{\rho(x, y)}\right) t}-\mathrm{e}^{-2 K\left(D_{\rho(x, y)}\right) T}}{1-\mathrm{e}^{-2 K\left(D_{\rho(x, y)}\right) T}} \mathrm{e}^{K\left(D_{\rho(x, y)}\right) t} \rho(x, y) \\
& \leqslant \rho(x, y) 1_{[0, T]}(t), \quad t \in\left[0, \Theta_{n}\right] .
\end{aligned}
$$

So, $\tau_{D(x, y)} \geqslant \tau_{D}(y)$ and $T \geqslant \tau$. Combining these inequalities with (2.5) we complete the proof.

### 2.2. Proof of (2) implying (3)

We will present below a more general result, which works for sub-Markovian operators on metric spaces. Let $(E, \rho)$ be a metric space, and let $P$ be a sub-Markovian operator on $\mathscr{B}_{b}(E)$.

$$
\delta(f)(x)=\limsup _{y \rightarrow x} \frac{|f(y)-f(x)|}{\rho(x, y)}, \quad x \in E, f \in \mathscr{B}_{b}(E) .
$$

If in particular $E=M$ and $f$ is differentiable at point $x$, then $\delta(f)(x)=|\nabla f|(x)$. So, (2) implying (3) is a direct consequence of the following result.

Proposition 2.3. Let $x \in E$ be fixed. If there exists a positive continuous function $\Phi$ on $E$ such that the log-Harnack inequality

$$
\begin{equation*}
P \log f(y) \leqslant \log \{P f(x)+1-P 1(x)\}+\Phi(y) \rho(x, y)^{2}, \quad f>0, f \in \mathscr{B}_{b}(E), \tag{2.6}
\end{equation*}
$$

holds for small $\rho(x, y)$, then

$$
\begin{equation*}
\delta(P f)^{2}(x) \leqslant 2 \Phi(x)\left\{P f^{2}(x)-(P f)^{2}(x)\right\}, \quad f \in \mathscr{B}_{b}(E) \tag{2.7}
\end{equation*}
$$

Proof. Let $f \in \mathscr{B}_{b}(E)$. According to the proof of [8, Proposition 2.3], (2.6) for small $\rho(x, y)$ implies that $\operatorname{Pf}$ is continuous at $x$. Let $\left\{x_{n}\right\}_{n} \geqslant 1$ be a sequence converging to $x$, and denote $\varepsilon_{n}=\rho\left(x_{n}, x\right)$. For any positive constant $c>0$, we apply (2.6) to $c \varepsilon_{n} f+1$ in place of $f$, so that for large enough $n$

$$
P \log \left(c \varepsilon_{n} f+1\right)\left(x_{n}\right) \leqslant \log \left\{P\left(c \varepsilon_{n} f+1\right)(x)+1-P 1(x)\right\}+\Phi\left(x_{n}\right) \varepsilon_{n}^{2} .
$$

Noting that for large $n$ (or for small $\varepsilon_{n}$ ) we have

$$
\begin{aligned}
& P \log \left(c \varepsilon_{n} f+1\right)\left(x_{n}\right)=P\left(c \varepsilon_{n} f-\frac{1}{2}\left(c \varepsilon_{n}\right)^{2} f^{2}\right)\left(x_{n}\right)+\mathrm{o}\left(\varepsilon_{n}^{2}\right) \\
& \\
& =c \varepsilon_{n} P f(x)+c \varepsilon_{n}^{2} \frac{P f\left(x_{n}\right)-P f(x)}{\rho\left(x_{n}, x\right)}-\frac{1}{2}\left(c \varepsilon_{n}\right)^{2} P f^{2}(x)+\mathrm{o}\left(\varepsilon_{n}^{2}\right),
\end{aligned} \quad \begin{aligned}
& \log \left\{P\left(c \varepsilon_{n} f+1\right)(x)+1-P 1(x)\right\}=c \varepsilon_{n} P f(x)-\frac{1}{2}\left(c \varepsilon_{n}\right)^{2}(P f)^{2}(x)+\mathrm{o}\left(\varepsilon_{n}^{2}\right) .
\end{aligned}
$$

We obtain

$$
c \limsup _{n \rightarrow \infty} \frac{P f\left(x_{n}\right)-P f(x)}{\rho\left(x_{n}, x\right)} \leqslant \frac{c^{2}}{2}\left\{P f^{2}(x)-(P f)^{2}(x)\right\}+\Phi(x), \quad c>0 .
$$

Exchanging the positions of $x_{n}$ and $x$, we also have

$$
c \limsup _{n \rightarrow \infty} \frac{P f(x)-P f\left(x_{n}\right)}{\rho\left(x_{n}, x\right)} \leqslant \frac{c^{2}}{2}\left\{P f^{2}(x)-(P f)^{2}(x)\right\}+\Phi(x), \quad c>0 .
$$

Therefore,

$$
\delta(P f)(x) \leqslant \frac{c}{2}\left\{P f^{2}(x)-(P f)^{2}(x)\right\}+\frac{\Phi(x)}{c}, \quad c>0 .
$$

This implies (2.7) by minimizing the upper bound in $c>0$.

### 2.3. Proof of (3) implying (1)

The proof of $\mathrm{Ric}_{Z} \geqslant-K$ is more or less standard by using the Taylor expansions for small $T>0$. Let $x \in M \backslash \partial M$ and $D=B(x, r) \subset M \backslash \partial M$ for small $r>0$ such that $\phi:=r^{2}-\rho(x, \cdot)^{2} \in$ $\mathscr{C}_{D}$. It is easy to see that for $f \in C_{0}^{\infty}(M)$ and small $t>0$,

$$
\begin{aligned}
& \left|\nabla P_{t} f\right|^{2}(x)=|\nabla f|^{2}(x)+2 t\langle\nabla f, \nabla L f\rangle+\mathrm{o}(t), \\
& \frac{K(D)}{1-\mathrm{e}^{-2 K(D) t}}+\frac{c_{D}(\phi)^{2}\left(\mathrm{e}^{2 K(D) t}-1\right)}{2 K(D) \phi(x)^{4}}=\frac{1}{2 t}+\frac{K(D)}{2}+\mathrm{o}(1) .
\end{aligned}
$$

Moreover (see [10, (3.6)]),

$$
P_{t} f^{2}(x)-\left(P_{t} f\right)^{2}(x)=2 t|\nabla f|^{2}(x)+t^{2}\left\{2\langle\nabla f, \nabla L f\rangle+L|\nabla f|^{2}\right\}(x)+\mathrm{o}(t)
$$

Combining these with (2.7) we obtain

$$
\Gamma_{2}(f)(x):=\frac{1}{2} L|\nabla f|^{2}(x)-\langle\nabla f, \nabla L f\rangle(x) \geqslant-K(D)|\nabla f|^{2}(x)=-\left(\sup _{B(x, r)} K\right)|\nabla f|^{2}(x) .
$$

Letting $r \downarrow 0$, we arrive at $\Gamma_{2}(f)(x) \geqslant-K(x)$ for $x \in M \backslash \partial M$ and $f \in C_{0}^{\infty}(M)$, which is equivalent to (1.1).

Next, we assume that $\partial M \neq \emptyset$ and intend to prove from (3) that the second fundamental form $\mathbb{I}$ of $\partial M$ is non-negative, i.e. $\partial M$ is convex. When $M$ is compact, the proof was done in [10] (see the proof of Theorem 1.1 therein for (7) implying (1)). Below we show that the proof works for general setting by using a localization argument with a stopping time.

Let $x \in \partial M$ and $r>0$. Define

$$
\sigma_{r}=\inf \left\{s \geqslant 0: \rho\left(X_{s}, x\right) \geqslant r\right\}
$$

where $X_{s}$ is the $L$-reflecting diffusion process starting at point $x$. Let $l_{s}$ be the local time of the process on $\partial M$. Then, according to [13, Lemmas 2.3 and 3.1], there exist two constants $C_{1}, C_{2}>0$ such that

$$
\begin{equation*}
\mathbb{P}\left(\sigma_{r} \leqslant t\right) \leqslant \mathrm{e}^{-C_{1} / t}, \quad t \in(0,1], \tag{2.8}
\end{equation*}
$$

and

$$
\begin{equation*}
\left|\mathbb{E} l_{t \wedge \sigma_{r}}-\frac{2 \sqrt{t}}{\sqrt{\pi}}\right| \leqslant C_{2} t, \quad t \in[0,1], \tag{2.9}
\end{equation*}
$$

where (2.8) is also ensured by [2, Lemma 2.3] for $\partial M=\emptyset$. Let $f \in C_{0}^{\infty}(M)$ satisfy the Neumann boundary condition. We aim to prove $\mathbb{I}(\nabla f, \nabla f)(x) \geqslant 0$. To apply Theorem 1.1(3), we construct $D$ and $\phi \in \mathscr{C}_{D}$ as follows.

Firstly, let $\varphi \in C_{0}^{\infty}(\partial M)$ such that $\varphi(x)=1$ and $\operatorname{supp} \varphi \subset \partial M \cap B(x, r / 2)$, where $B(x, s)=$ $\{z \in M: \rho(z, x)<s\}$ for $s>0$. Then letting $\phi_{0}\left(\exp _{y}[s N]\right)=\varphi(y)$ (where $y \in \partial M, s \geqslant 0$ ), we extend $\varphi$ to a smooth function in a neighborhood of $\partial M$, say $\partial_{r_{0}} M:=\left\{z \in M: \rho(z, \partial M)<r_{0}\right\}$ for some $r_{0} \in(0, r)$ such that $\rho(\cdot, \partial M)$ is smooth on $\left(\partial_{r_{0}} M\right) \cap B(x, r)$. Obviously, $\phi_{0}$ satisfies the Neumann boundary condition. Finally, for $h \in C^{\infty}([0, \infty))$ with $\left.h\right|_{\left[0, r_{0} / 4\right]}=1$ and $\left.h\right|_{\left[r_{0} / 2, \infty\right)}=0$, we take $\phi=\phi_{0} h(\rho(\cdot, \partial M))$ and $D=\{z \in M: \phi(z)>0\}$. Then $\phi(x)=1$, $\left.\phi\right|_{\partial D \backslash \partial M}=0,\left.N \phi\right|_{\partial M}=\left.N \phi_{0}\right|_{\partial M}=0$, and $D \subset B(x, r)$.

Once $D$ and $\phi \in \mathscr{C}_{D}$ are given, below we calculate both sides of the gradient inequality in (3) respectively.

According to (2.8), for small $t>0$ we have

$$
\begin{align*}
P_{t} f^{2}(x) & =\mathbb{E} f^{2}\left(X_{t \wedge \sigma_{r}}\right)+\mathrm{o}\left(t^{2}\right)=f^{2}(x)+\mathbb{E} \int_{0}^{t \wedge \sigma_{r}} L f^{2}\left(X_{s}\right) \mathrm{d} s+\mathrm{o}\left(t^{2}\right) \\
& =f^{2}(x)+2 \mathbb{E} \int_{0}^{t \wedge \sigma_{r}}(f L f)\left(X_{s}\right) \mathrm{d} s+2 \mathbb{E} \int_{0}^{t \wedge \sigma_{r}}|\nabla f|^{2}\left(X_{s}\right) \mathrm{d} s+\mathrm{o}\left(t^{2}\right) \tag{2.10}
\end{align*}
$$

Noting that by the Neumann boundary condition

$$
\mathbb{E}\left|f(x)-f\left(X_{s \wedge \sigma_{r}}\right)\right|^{2} \leqslant\left\|L(f(x)-f)^{2}\right\|_{\infty} s, \quad s \geqslant 0
$$

we have

$$
\begin{align*}
& \mathbb{E} \int_{0}^{t \wedge \sigma_{r}}(f L f)\left(X_{s}\right) \mathrm{d} s-f(x) \mathbb{E} \int_{0}^{t \wedge \sigma_{r}} L f\left(X_{s}\right) \mathrm{d} s \\
& \quad=\mathbb{E} \int_{0}^{t \wedge \sigma_{r}} L f(x)\left\{f(x)-f\left(X_{s}\right)\right\} \mathrm{d} s+\mathbb{E} \int_{0}^{t \wedge \sigma_{r}}\left(L f\left(X_{s}\right)-L f(x)\right)\left(f(x)-f\left(X_{s}\right)\right) \mathrm{d} s \\
& \quad \leqslant\|L f\|_{\infty} \mathbb{E} \int_{0}^{t \wedge \sigma_{r}} \int_{0}^{s} \mathrm{~d} u+\mathbb{E} \int_{0}^{t} \sqrt{\mathbb{E}\left|L f\left(X_{s \wedge \sigma_{r}}\right)-L f(x)\right|^{2} \cdot \mathbb{E}\left|f(x)-f\left(X_{s \wedge \sigma_{r}}\right)\right|^{2}} \mathrm{~d} s \\
& \quad=\mathrm{o}\left(t^{3 / 2}\right) . \tag{2.11}
\end{align*}
$$

Moreover, by the Itô formula and the fact that $N|\nabla f|^{2}=2 \mathbb{I}(\nabla f, \nabla f)$ holds on $\partial M$ (see e.g. [10, (3.8)]), we have

$$
\begin{aligned}
\mathbb{E}|\nabla f|^{2}\left(X_{s \wedge \sigma_{r}}\right) & =|\nabla f|^{2}(x)+\mathbb{E} \int_{0}^{s \wedge \sigma_{r}} L|\nabla f|^{2}\left(X_{u}\right) \mathrm{d} u+2 \int_{0}^{s \wedge \sigma_{r}} \mathbb{I}(\nabla f, \nabla f)\left(X_{u}\right) \mathrm{d} l_{u} \\
& \leqslant|\nabla f|^{2}(x)+2 \mathbb{I}(r) \mathbb{E} l_{s \wedge \sigma_{r}}+\mathrm{O}(t),
\end{aligned}
$$

where

$$
\mathbb{I}(r):=\sup \{\mathbb{I}(\nabla f, \nabla f)(y): y \in \partial M \cap B(x, r)\}
$$

Combining this with (2.10), (2.11) and using (2.8) and (2.9), we obtain

$$
\begin{equation*}
P_{t} f^{2}(x) \leqslant f^{2}(x)+2 f(x) \mathbb{E} \int_{0}^{t \wedge \sigma_{r}} L f\left(X_{s}\right) \mathrm{d} s+C t^{3 / 2} \mathbb{I}(r)+\mathrm{o}\left(t^{3 / 2}\right) \tag{2.12}
\end{equation*}
$$

for some constant $C>0$ and small $t>0$.
On the other hand, by (2.8) we have

$$
\begin{aligned}
\left(P_{t} f\right)^{2}(x) & =\left(f(x)+\mathbb{E} \int_{0}^{t \wedge \sigma_{r}} L f\left(X_{s}\right) \mathrm{d} s+\mathrm{o}\left(t^{2}\right)\right)^{2} \\
& =f^{2}(x)+2 t f(x) \mathbb{E} \int_{0}^{t \wedge \sigma_{r}} L f\left(X_{s}\right) \mathrm{d} s+\mathrm{o}\left(t^{2}\right) .
\end{aligned}
$$

Combining this with (2.12) and noting that

$$
\frac{K(D)}{1-\mathrm{e}^{-2 K(D) t}}+\frac{c_{D}(\phi)^{2}\left(\mathrm{e}^{2 K(D) t}-1\right)}{2 K(D) \phi(x)^{4}}=\frac{1}{2 t}+\mathrm{O}(1)
$$

holds for small $t>0$, we arrive at

$$
\begin{aligned}
& \left\{P_{t} f^{2}-\left(P_{t} f\right)^{2}\right\}(x)\left(\frac{K(D)}{1-\mathrm{e}^{-2 K(D) t}}+\frac{c_{D}(\phi)^{2}\left(\mathrm{e}^{2 K(D) t}-1\right)}{2 K(D) \phi(x)^{4}}\right) \\
& \quad \leqslant|\nabla f|^{2}(x)+C \mathbb{I}(r) \sqrt{t}+\mathrm{o}\left(t^{1 / 2}\right)
\end{aligned}
$$

for small $t>0$. Combining this with the gradient inequality in (3) and noting that

$$
\left|\nabla P_{t} f\right|^{2}(x)=\left|\nabla f(x)+\int_{0}^{t} \nabla P_{s} L f(x) \mathrm{d} s\right|^{2}=|\nabla f|^{2}(x)+\mathrm{O}(t)
$$

we conclude that

$$
\begin{aligned}
\mathbb{I}(r) \geqslant & \lim _{t \rightarrow 0} \frac{1}{C \sqrt{t}}\left\{\left\{P_{t} f^{2}-\left(P_{t} f\right)^{2}\right\}(x)\left(\frac{K(D)}{1-\mathrm{e}^{-2 K(D) t}}+\frac{c_{D}(\phi)^{2}\left(\mathrm{e}^{2 K(D) t}-1\right)}{2 K(D) \phi(x)^{4}}\right)\right. \\
& \left.-\left|\nabla P_{t} f\right|^{2}(x)\right\} \geqslant 0 .
\end{aligned}
$$

Therefore, $\mathbb{I}(\nabla f, \nabla f)(x)=\lim _{r \rightarrow 0} \mathbb{I}(r) \geqslant 0$.

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